

Math 6810
(Probability and Fractals)

Spring 2016

Lecture notes

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Recommended reading: (Do not purchase these books before consulting with your instructor!)

1. *Real Analysis* by H. L. Royden (4th edition), Prentice Hall.
2. *Probability and Measure* by P. Billingsley (3rd edition), Wiley.
3. *Probability with Martingales* by D. Williams, Cambridge University Press.
4. *Fractal Geometry: Foundations and Applications* by K. Falconer (2nd edition), Wiley.

Chapter 4

Random fractals

This chapter is based on Falconer, Chapters 4 and 15, and a paper by B. Hunt.

4.1 The potential theoretic method

Proposition 4.1. *Let μ be a mass distribution on \mathbb{R}^n , $F \subset \mathbb{R}^n$ a Borel set, and $0 < c < \infty$ a constant. Suppose that*

$$\limsup_{r \downarrow 0} \frac{\mu(B(x, r))}{r^s} < c \quad \text{for all } x \in F. \quad (4.1)$$

Then $\mathcal{H}^s(F) \geq \mu(F)/c$.

Proof. For $\delta > 0$, let

$$F_\delta := \{x \in F : \mu(B(x, r)) < cr^s \text{ for all } 0 < r \leq \delta\}.$$

Let $\{U_i\}$ be a δ -cover of F . If $x \in U_i \cap F_\delta$, then $U_i \subset B(x, |U_i|)$ and

$$\mu(U_i) \leq \mu(B(x, |U_i|)) < c|U_i|^s.$$

It follows that

$$\mu(F_\delta) \leq \sum_i \{\mu(U_i) : U_i \cap F_\delta \neq \emptyset\} \leq c \sum_i |U_i|^s.$$

Taking the infimum over all δ -covers gives $\mu(F_\delta) \leq c\mathcal{H}_\delta^s(F)$. But, since F_δ increases to F as $\delta \rightarrow 0$ by (4.1), we have $\mu(F) = \lim_{\delta \rightarrow 0} \mu(F_\delta)$, and so $\mu(F) \leq c\mathcal{H}^s(F)$. \square

Definition 4.2. Let μ be a mass distribution on \mathbb{R}^n and $s \geq 0$. The s -potential at a point $x \in \mathbb{R}^n$ due to μ is

$$\phi_s(x) := \int \frac{d\mu(y)}{|x - y|^s}, \quad (4.2)$$

and the s -energy of μ is

$$I_s(\mu) := \int \phi_s(x) d\mu(x) = \iint \frac{d\mu(x) d\mu(y)}{|x - y|^s}, \quad (4.3)$$

where each integral is taken over \mathbb{R}^n .

Observe that if $I_s(\mu) < \infty$, then μ is nonatomic; that is, $\mu(\{x\}) = 0$ for each $x \in \mathbb{R}^n$. The connection between s -energy and Hausdorff measure and Hausdorff dimension is as follows:

Theorem 4.3. *Let $F \subset \mathbb{R}^n$. If there is a mass distribution μ on F with $I_s(\mu) < \infty$, then $\mathcal{H}^s(F) = \infty$ and therefore, $\dim_H F \geq s$.*

(There is a sort-of-converse to this theorem, which we will not need - see Falconer, Theorem 4.13(b).)

Proof. Let μ be a mass distribution on F , so $\mu(F) > 0 = \mu(\mathbb{R}^n \setminus F)$, and suppose $I_s(\mu) < \infty$. Define the set

$$F_1 := \left\{ x \in F : \limsup_{r \downarrow 0} \frac{\mu(B(x, r))}{r^s} > 0 \right\}.$$

Fix $x \in F_1$. Then we can find $\varepsilon > 0$ and a sequence $\{r_i\}$ decreasing to 0 such that $\mu(B(x, r_i)) \geq \varepsilon r_i^s$ for each i . Since $\mu(\{x\}) = 0$, we can find $0 < q_i < r_i$ small enough so that $\mu(A_i) \geq \frac{1}{2}\varepsilon r_i^s$, where A_i is the annulus $A_i = B(x, r_i) \setminus B(x, q_i)$. Taking a subsequence if necessary, we may assume $r_{i+1} < q_i$, so that the annuli A_i are disjoint. It follows that

$$\phi_s(x) = \int \frac{d\mu(y)}{|x-y|^s} \geq \sum_{i=1}^{\infty} \int_{A_i} \frac{d\mu(y)}{|x-y|^s} \geq \sum_{i=1}^{\infty} \frac{1}{2} \varepsilon r_i^s r_i^{-s} = \infty.$$

On the other hand, $I_s(\mu) = \int \phi_s(x) d\mu(x) < \infty$, so $\phi_s(x) < \infty$ for μ -almost every x . Though it's not clear if F_1 is measurable (a Borel set), there is certainly a Borel set E such that $F_1 \subset E$ and $\mu(E) = 0$, and by definition of F_1 ,

$$\limsup_{r \downarrow 0} \frac{\mu(B(x, r))}{r^s} = 0 \quad \text{for all } x \in F \setminus E.$$

Hence, by Proposition 4.1,

$$\mathcal{H}^s(F) \geq \mathcal{H}^s(F \setminus E) \geq \mu(F \setminus E)/c = \mu(F)/c,$$

for every $c > 0$. Therefore, $\mathcal{H}^s(F) = \infty$. □

We can use Theorem 4.3 to obtain (almost-sure) lower bounds for the Hausdorff dimension of random fractals as follows. Let (Ω, \mathcal{F}, P) be a probability space, and suppose for each $\omega \in \Omega$ we have a fractal F_ω . If we can find for each ω a mass distribution μ_ω on F_ω such that

$$\int_{\Omega} I_s(\mu_\omega) dP(\omega) = \int_{\Omega} \iint \frac{d\mu_\omega(x) d\mu_\omega(y)}{|x-y|^s} dP(\omega) < \infty,$$

then it will follow that $I_s(\mu_\omega) < \infty$ for almost every ω , and so $\dim_H F_\omega \geq s$ for almost every ω , in other words, with probability one. In many practical applications a suitable change of variable and Fubini's theorem can be applied to compute or estimate the above triple integral.

4.2 Random Cantor sets

We will construct a “statistically self-similar” set analogous to the Cantor set by randomly choosing the contraction ratios at each stage of the construction. Let $0 < a \leq b < \frac{1}{2}$ be constants. Let $(\Omega, \mathcal{F}, \mathbb{P})$ be a probability space on which is defined a collection of random variables C_{i_1, \dots, i_k} , $k \in \mathbb{N}$, $(i_1, \dots, i_k) \in \mathcal{I}_k := \{1, 2\}^k$, with the following properties:

- (i) For each $k \in \mathbb{N}$ and $(i_1, \dots, i_k) \in \mathcal{I}_k$, the random variables $C_{i_1, \dots, i_k, j}$, $j = 1, 2$ are independent of \mathcal{F}_k , where

$$\mathcal{F}_k := \sigma(\{C_{i_1, \dots, i_l} : 1 \leq l \leq k, (i_1, \dots, i_l) \in \mathcal{I}_l\}).$$

- (ii) For each $k \in \mathbb{N}$ the collection of random pairs $\{(C_{i_1, \dots, i_k, 1}, C_{i_1, \dots, i_k, 2}) : (i_1, \dots, i_k) \in \mathcal{I}_k\}$ is independent.
- (iii) For each $k \in \mathbb{N}$ and $(i_1, \dots, i_k) \in \mathcal{I}_k$, $C_{i_1, \dots, i_k, 1} \sim C_1$ and $C_{i_1, \dots, i_k, 2} \sim C_2$.
- (iv) $a \leq C_j \leq b$ a.s. for $j = 1, 2$.

Note that we do not require independence of $C_{i_1, \dots, i_k, 1}$ and $C_{i_1, \dots, i_k, 2}$.

Now define a collection of random intervals $\{I_{i_1, \dots, i_k}\}$ as follows: put $I_1 = [0, C_1]$ and $I_2 = [1 - C_2, 1]$. For $k \in \mathbb{N}$ and $(i_1, \dots, i_k) \in \mathcal{I}_k$, let $I_{i_1, \dots, i_k, 1}$ and $I_{i_1, \dots, i_k, 2}$ be subintervals of $I := I_{i_1, \dots, i_k}$ such that $I_{i_1, \dots, i_k, 1}$ has the same left endpoint as I , $I_{i_1, \dots, i_k, 2}$ has the same right endpoint as I , and $|I_{i_1, \dots, i_k, j}| = C_{i_1, \dots, i_k, j}|I|$ for $j = 1, 2$. We call I_{i_1, \dots, i_k} a *basic interval at level k* . Let

$$E_k = \bigcup_{(i_1, \dots, i_k) \in \mathcal{I}_k} I_{i_1, \dots, i_k}, \quad k \in \mathbb{N},$$

and

$$F := \bigcap_{k=1}^{\infty} E_k. \quad (4.4)$$

Note that for each k , the basic intervals of level k are disjoint, and as a result, F has the topological properties of a Cantor set (perfect, totally disconnected).

Theorem 4.4. *For the random Cantor set F described above, $\dim_H F = s$ with probability 1, where s is the unique positive solution of*

$$\mathbb{E}(C_1^s + C_2^s) = 1. \quad (4.5)$$

Proof. It is straightforward to check (using that C_1 and C_2 are bounded random variables) that $f(s) := \mathbb{E}(C_1^s + C_2^s)$ is continuous and strictly decreasing in s , with $f(0) = 2$ and $f(1) = \mathbb{E}(C_1 + C_2) \leq 2b < 1$, so (4.5) has a unique solution.

Let \mathcal{E}_k be the (finite) collection of intervals I_{i_1, \dots, i_k} , $(i_1, \dots, i_k) \in \mathcal{I}_k$, with $\mathcal{E}_0 := \{[0, 1]\}$. For $I = I_{i_1, \dots, i_k} \in \mathcal{E}_k$, write $I_L := I_{i_1, \dots, i_k, 1}$ and $I_R := I_{i_1, \dots, i_k, 2}$. For $s > 0$, we have

$$\begin{aligned} \mathbb{E}(|I_{i_1, \dots, i_k, 1}|^s + |I_{i_1, \dots, i_k, 2}|^s | \mathcal{F}_k) &= \mathbb{E}(C_{i_1, \dots, i_k, 1}^s + C_{i_1, \dots, i_k, 2}^s) |I_{i_1, \dots, i_k}|^s \\ &= \mathbb{E}(C_1^s + C_2^s) |I_{i_1, \dots, i_k}|^s, \end{aligned}$$

where we first used (i) and then (iii). Summing over all the intervals in \mathcal{E}_k gives

$$\mathbb{E} \left(\sum_{I \in \mathcal{E}_{k+1}} |I|^s \middle| \mathcal{F}_k \right) = \sum_{I \in \mathcal{E}_k} |I|^s \mathbb{E}(C_1^s + C_2^s). \quad (4.6)$$

Taking expectations on both sides we obtain

$$\mathbb{E} \left(\sum_{I \in \mathcal{E}_{k+1}} |I|^s \right) = \mathbb{E} \left(\sum_{I \in \mathcal{E}_k} |I|^s \right) \mathbb{E}(C_1^s + C_2^s). \quad (4.7)$$

If s is the solution of (4.5), then (4.6) reduces to

$$\mathbb{E} \left(\sum_{I \in \mathcal{E}_{k+1}} |I|^s \middle| \mathcal{F}_k \right) = \sum_{I \in \mathcal{E}_k} |I|^s,$$

which shows that the sequence (X_k) defined by $X_k = \sum_{I \in \mathcal{E}_k} |I|^s$ is a martingale.

Exercise: Show that the martingale (X_k) is L^2 -bounded. (*Hint:* Show that $\mathbb{E}(X_{k+1}^2 | \mathcal{F}_k) \leq X_k^2 + a\gamma^k$, where $\gamma = \mathbb{E}(C_1^{2s} + C_2^{2s}) < 1$ and a is a constant.)

As a result, $X := \lim_{k \rightarrow \infty} X_k$ exists almost surely, and $\mathbb{E}(X) = \mathbb{E}(X_0) = 1$. We claim that $X > 0$ almost surely. Let $q = \mathbb{P}(X = 0)$. Since $X \geq 0$ and $\mathbb{E}(X) = 1$, $q < 1$. Now

$$X_k = \sum_{I \in \mathcal{E}_k, I \subset I_1} |I|^s + \sum_{I \in \mathcal{E}_k, I \subset I_2} |I|^s,$$

and the two random sums on the right are independent by (ii) (for all $k \geq 2$), and each tends to 0 with probability q , by the self-similarity of the construction. Thus $q = \mathbb{P}(X_k \rightarrow 0) = q^2$, and so $q = 0$, proving the claim. It follows that there are random variables M_1 and M_2 such that

$$0 < M_1 \leq X_k = \sum_{I \in \mathcal{E}_k} |I|^s \leq M_2 < \infty \quad \text{a.s. for all } k. \quad (4.8)$$

Given $\delta > 0$, \mathcal{E}_k is a δ -cover of F for large enough k , and so $\mathcal{H}^s(F) \leq M_2 < \infty$ with probability 1. Hence, $\dim_H F \leq s$ almost surely.

For the lower bound we use the potential theoretic method. Let s again be the solution of (4.5). For $I \in \mathcal{E}_k$, define the random variable

$$\mu(I) := \lim_{j \rightarrow \infty} \sum \{ |J|^s : J \in \mathcal{E}_j, J \subset I \}. \quad (4.9)$$

By the same argument as above, this limit exists, is \mathcal{F}_k -measurable, and $0 < \mu(I) < \infty$ almost surely. Furthermore, if $I \in \mathcal{E}_k$, then $\mu(I) = \mu(I_L) + \mu(I_R)$, so μ extends to a (random!) mass distribution on $[0, 1]$ with support in F . (The complete proof of this fact is rather involved and is omitted here.) In addition, we have

$$\mathbb{E}[\mu(I) | \mathcal{F}_k] = |I|^s, \quad I \in \mathcal{E}_k. \quad (4.10)$$

Fix $0 < t < s$. We will estimate the t -energy of μ . For $x, y \in F$, there is a largest k such that x and y belong to the same basic interval $I \in \mathcal{E}_k$; denote this interval by $x \wedge y$. If $I \in \mathcal{E}_k$, then the subintervals I_L and I_R are separated by a gap of length at least $d|I|$, where $d = 1 - 2b > 0$. Thus, for any $I \in \mathcal{E}_k$,

$$\begin{aligned} \iint_{x \wedge y = I} |x - y|^{-t} d\mu(x) d\mu(y) &= 2 \int_{I_L} \int_{I_R} |x - y|^{-t} d\mu(x) d\mu(y) \\ &\leq 2d^{-t} |I|^{-t} \mu(I_L) \mu(I_R), \end{aligned}$$

and so

$$\begin{aligned} \mathbb{E} \left(\iint_{x \wedge y = I} |x - y|^{-t} d\mu(x) d\mu(y) \middle| \mathcal{F}_{k+1} \right) &\leq 2d^{-t} |I|^{-t} \mathbb{E}[\mu(I_L) | \mathcal{F}_{k+1}] \mathbb{E}[\mu(I_R) | \mathcal{F}_{k+1}] \\ &\leq 2d^{-t} |I|^{-t} |I_L|^s |I_R|^s \\ &\leq 2d^{-t} |I|^{2s-t}. \end{aligned}$$

Here the first inequality uses that, conditionally on \mathcal{F}_{k+1} , $\mu(I_R)$ and $\mu(I_L)$ are independent because of assumption (ii); and the second inequality follows from (4.10). Taking expectations we obtain

$$\mathbb{E} \left(\iint_{x \wedge y = I} |x - y|^{-t} d\mu(x) d\mu(y) \right) \leq 2d^{-t} \mathbb{E}(|I|^{2s-t}).$$

Summing over $I \in \mathcal{E}_k$ and iterating (4.7), we get

$$\mathbb{E} \left(\sum_{I \in \mathcal{E}_k} \iint_{x \wedge y = I} |x - y|^{-t} d\mu(x) d\mu(y) \right) \leq 2d^{-t} \mathbb{E} \left(\sum_{I \in \mathcal{E}_k} |I|^{2s-t} \right) = 2d^{-t} \lambda^k,$$

where $\lambda := \mathbb{E}(C_1^{2s-t} + C_2^{2s-t}) < 1$, since $2s - t > s$. Finally, we can sum over k to obtain

$$\begin{aligned} \mathbb{E} \left(\iint_F |x - y|^{-t} d\mu(x) d\mu(y) \right) &= \mathbb{E} \left(\sum_{k=0}^{\infty} \sum_{I \in \mathcal{E}_k} \iint_{x \wedge y = I} |x - y|^{-t} d\mu(x) d\mu(y) \right) \\ &\leq \sum_{k=0}^{\infty} 2d^{-t} \lambda^k < \infty. \end{aligned}$$

This implies that the t -energy of μ is finite almost surely, and hence, by Theorem 4.3, $\dim_H F \geq t$ a.s. Since $t < s$ was arbitrary, it follows that $\dim_H F \geq s$ almost surely. \square

Remark 4.5. Note that the proof does not tell us whether $\mathcal{H}^s(F) > 0$. The condition that there is a minimum gap between basic intervals is not necessary. (A version of the open set condition is enough.) But without this assumption, the proof is more involved.

Example 4.6. Let U be a uniformly distributed random variable on $(1/3, 2/3)$. Consider the construction of a random Cantor set F whereby for each basic interval $I = I_{i_1, \dots, i_k}$, the middle portion of length $U_{i_1, \dots, i_k} |I|$ is removed from I , where the collection $\{U_{i_1, \dots, i_k} : k \in \mathbb{N}\}$

$\mathbb{N}_0, (i_1, \dots, i_k) \in \mathcal{I}_k\}$ is independent and each $U_{i_1, \dots, i_k} \sim U$. This fits the framework of the above theorem, with $C_1 = C_2 = (1 - U)/2$. Thus, $\dim_H F = s$, where s is the solution of

$$\mathbb{E}(C_1^s + C_2^s) = 2 \mathbb{E} \left[\left(\frac{1 - U}{2} \right)^s \right] = 1. \quad (4.11)$$

Exercise 4.7. Solve (4.11) (numerically). (Solution: $\dim_H F = s \doteq .4966$).

4.3 A random Weierstrass function

In this section we randomize the construction of the Weierstrass function from Chapter 3 by adding random phases as follows. Let $\theta_1, \theta_2, \dots$ be independent uniform $(0, 2\pi)$ random variables. For constants $\lambda > 1$ and $1 < s < 2$, define the random function

$$W(x) = \sum_{n=0}^{\infty} \lambda^{(s-2)n} \sin(\lambda^n x + \theta_n), \quad 0 \leq x \leq 1. \quad (4.12)$$

The following theorem is due to B. Hunt (“The Hausdorff dimension of graphs of Weierstrass functions”, *Proc. Amer. Math. Soc.* 126 (1998), no. 3, 791–800). We present his proof with minor changes in notation.

Theorem 4.8. *With probability one, $\dim_H \text{Graph}(W) = s$.*

The proof uses convolutions of densities. We need a definition and some lemmas.

Definition 4.9. The *convolution* of two functions f and g in $L^1(\mathbb{R})$ is the function

$$f * g(x) := \int_{\mathbb{R}} f(y)g(x - y)dy.$$

An easy exercise (using Fubini’s theorem) shows that $f * g$ is well defined and in $L^1(\mathbb{R})$.

Lemma 4.10. *Let X and Y be independent random variables and suppose X is absolutely continuous with density f_X . Then $X + Y$ is absolutely continuous with density f_{X+Y} , and $\sup_z f_{X+Y}(z) \leq \sup_x f_X(x)$.*

Proof. Let μ_Y denote the distribution of Y , F_X the c.d.f. of X , and F_{X+Y} the c.d.f. of $X + Y$. Then, by integrating over the half-plane $x + y \leq z$ and using Fubini’s theorem,

$$F_{X+Y}(z) = \int_{\mathbb{R}} F_X(z - y)d\mu_Y(y), \quad z \in \mathbb{R}.$$

Since F_X is absolutely continuous, it now follows easily that F_{X+Y} is absolutely continuous also. (Check this!) Then by differentiating both sides of the above equation,

$$f_{X+Y}(z) = \int_{\mathbb{R}} f_X(z - y)d\mu_Y(y),$$

from which the second statement of the lemma follows immediately. \square

Lemma 4.11. *If X and Y are independent absolutely continuous random variables with densities f_X and f_Y , then $X + Y$ has density $f_{X+Y} = f_X * f_Y$.*

Proof. Easy exercise. □

Lemma 4.12 (Young's inequality for convolutions). *Let p, q and r be real numbers in $(1, \infty)$ such that $\frac{1}{p} + \frac{1}{q} = \frac{1}{r} + 1$. If $f \in L^p$ and $g \in L^q$, then $f * g \in L^r$, and*

$$\|f * g\|_r \leq \|f\|_p \|g\|_q.$$

(For a proof, which uses a generalized Hölder inequality, see proofwiki.org.)

Exercise 4.13. Let $q \neq 0$, $a \in \mathbb{R}$, and let θ be a uniform $(0, 2\pi)$ random variable. Show that the random variable $X = q \cos(a + \theta)$ has density

$$f_X(x) = \begin{cases} \frac{1}{\pi\sqrt{q^2-x^2}} & \text{if } |x| < |q|, \\ 0 & \text{if } |x| \geq |q|. \end{cases}$$

Finally, we recall the trigonometric identity

$$\sin x - \sin y = 2 \cos\left(\frac{x+y}{2}\right) \sin\left(\frac{x-y}{2}\right). \quad (4.13)$$

Proof of Theorem 4.8. First, it follows just as in the proof of Theorem 3.42 that W is Hölder continuous with exponent $2 - s$, and hence, by Proposition 3.41, $\dim_H \text{Graph}(W) \leq s$.

For the lower bound we use the potential-theoretic method. Let μ be the random measure on \mathbb{R}^2 , supported on $\text{Graph}(W)$, defined by

$$\mu(A) := \mathcal{L}(\{x \in [0, 1] : (x, W(x)) \in A\}), \quad A \in \mathcal{B}(\mathbb{R}^2),$$

where \mathcal{L} denotes Lebesgue measure on $[0, 1]$. If A_1, A_2, \dots are disjoint subsets of \mathbb{R}^2 , then the sets $\{x : (x, W(x)) \in A_i\}$ are disjoint, so μ is indeed a measure, and $\mu(\text{Graph}(W)) = 1$.

Fix $1 < t < s$. Our goal is to show that $E(I_t(\mu)) < \infty$, which will imply, as in the proof of Theorem 4.4, that $\dim_H \text{Graph}(W) \geq t$ a.s., and so, since t is arbitrary, $\dim_H \text{Graph}(W) \geq s$ a.s. Here we let \mathbf{x}, \mathbf{y} denote points in \mathbb{R}^2 . By a change-of-variable and the Pythagorean theorem,

$$I_t(\mu) = \int_{\mathbb{R}^2} \int_{\mathbb{R}^2} \frac{d\mu(\mathbf{x})d\mu(\mathbf{y})}{|\mathbf{x} - \mathbf{y}|^t} = \int_0^1 \int_0^1 \frac{dx dy}{((x-y)^2 + (W(x) - W(y))^2)^{t/2}},$$

so by Fubini's theorem,

$$E(I_t(\mu)) = \int_0^1 \int_0^1 E \left[((x-y)^2 + (W(x) - W(y))^2)^{-t/2} \right] dx dy. \quad (4.14)$$

Let $E_{x,y}$ denote the expectation in the above double integral. We will estimate $E_{x,y}$ for all $x, y \in [0, 1]$. Note first that, if $|x - y| \geq \pi/\lambda^2$, then $E_{x,y} \leq (\lambda^2/\pi)^t$. Fix now x and y with

$0 < |x - y| < \pi/\lambda^2$, and let $Z = W(x) - W(y)$. Then Z is a random variable, and we wish to show that Z has a bounded density function $h(z)$ satisfying

$$h(z) \leq C|x - y|^{s-2}, \quad z \in \mathbb{R}, \quad (4.15)$$

for some constant $C > 0$ that is independent of x and y . If we can show this, it will follow that

$$\begin{aligned} \mathbb{E}_{x,y} &= \mathbb{E} \left[((x - y)^2 + Z^2)^{-t/2} \right] = \int_{-\infty}^{\infty} \frac{h(z) dz}{((x - y)^2 + z^2)^{t/2}} \\ &\leq C \int_{-\infty}^{\infty} \frac{|x - y|^{s-2} dz}{((x - y)^2 + z^2)^{t/2}} \\ &= C \int_{-\infty}^{\infty} \frac{|x - y|^{s-2} |x - y| dw}{|x - y|^t (1 + w^2)^{t/2}} \\ &= C|x - y|^{s-1-t} \int_{-\infty}^{\infty} \frac{dw}{(1 + w^2)^{t/2}}, \end{aligned}$$

where the next-to-last step uses the change-of-variable $z = |x - y|w$. Since $t > 1$, the integral in the last line above converges. And (check!) since $t < s$, the double integral

$$\int_0^1 \int_0^1 |x - y|^{s-1-t} dx dy$$

converges also. Thus, by (4.14), $\mathbb{E}(I_t(\mu)) < \infty$, as desired.

Using (4.13), we can write

$$\begin{aligned} Z &= \sum_{n=0}^{\infty} \lambda^{(s-2)n} (\sin(\lambda^n x + \theta_n) - \sin(\lambda^n y + \theta_n)) \\ &= \sum_{n=0}^{\infty} 2\lambda^{(s-2)n} \sin\left(\lambda^n \frac{x - y}{2}\right) \cos\left(\lambda^n \frac{x + y}{2} + \theta_n\right) \\ &= \sum_{n=0}^{\infty} q_n \cos(r_n + \theta_n) =: \sum_{n=0}^{\infty} Z_n. \end{aligned}$$

Note that the random variables Z_1, Z_2, \dots are independent, with Z_n having density

$$h_n(z) = \begin{cases} \frac{1}{\pi\sqrt{q^2 - z^2}} & \text{if } |z| < |q|, \\ 0 & \text{if } |z| \geq |q|, \end{cases}$$

by Exercise 4.13. It follows from Lemma 4.10 that Z is absolutely continuous with density $h = h_0 * h_1 * h_2 \cdots$, and h is bounded by any upper bound for any finite convolution $h_j * \cdots * h_k$, where $j \leq k$.

Next, recall that $|x - y| < \pi/\lambda^2$. Thus, there is an integer $k \geq 2$ such that $\pi\lambda^{-k-1} < |x - y| \leq \pi\lambda^{-k}$. Fix this k . Then

$$\frac{\pi}{2\lambda^3} < \left| \lambda^{k-2} \frac{x - y}{2} \right| < \left| \lambda^k \frac{x - y}{2} \right| \leq \frac{\pi}{2},$$

and hence,

$$|q_n| > 2 \sin\left(\frac{\pi}{2\lambda^3}\right) \lambda^{(s-2)k} > 2 \sin\left(\frac{\pi}{2\lambda^3}\right) |x-y|^{2-s} \quad (4.16)$$

for $n = k-2, k-1, k$. Observe that $h_n \in L^p$ for $p < 2$, and by direct calculation, for $n = k-2, k-1, k$,

$$\|h_n\|_{3/2}^{3/2} = |q_n|^{-1/2} \int_{-1}^1 \frac{dw}{\pi(1-w^2)^{3/4}} = K|q_n|^{-1/2},$$

so that $\|h_n\|_{3/2} = K^{2/3}|q_n|^{-1/3} \leq K'|x-y|^{(s-2)/3}$ by (4.16), where K' depends only on λ . Now we apply first Young's inequality (Lemma 4.12) to obtain

$$\|h_{k-1} * h_k\|_3 \leq \|h_{k-1}\|_{3/2} \|h_k\|_{3/2},$$

and then Hölder's inequality to conclude that

$$\begin{aligned} h_{k-2} * h_{k-1} * h_k(z) &= \int_{\mathbb{R}} h_{k-2}(z-x)(h_{k-1} * h_k)(x) dx \\ &\leq \|h_{k-2}\|_{3/2} \|h_{k-1} * h_k\|_3 \\ &\leq \|h_{k-2}\|_{3/2} \|h_{k-1}\|_{3/2} \|h_k\|_{3/2} \\ &\leq K'^3 |x-y|^{s-2}. \end{aligned}$$

But then this bound also applies to $h(z)$, and so we arrive at (4.15), completing the proof. \square